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2006 J. Phys. A: Math. Gen. 39 2325

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Determinantal formulae for the Casimir operators of inhomogeneous Lie algebras

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Received 30 June 2005, in final form 25 January 2006

Published 22 February 2006

Online at stacks.iop.org/JPhysA/39/2325

Abstract

Contractions of Lie algebras are combined with the classical matrix method of Gel'fand to obtain matrix formulae for the Casimir operators of inhomogeneous Lie algebras. The method is presented for the inhomogeneous pseudo-unitary Lie algebras $Iu(p, q)$. This procedure is extended to contractions of $Iu(p, q)$ isomorphic to an extension by a derivation of the inhomogeneous special pseudo-unitary Lie algebras $I\mathfrak{su}(p-1, q)$, providing an additional analytical method to obtain their invariants. Further, matrix formulae for the invariants of other inhomogeneous Lie algebras are presented.

PACS number: 02.20.Sv

1. Introduction

Unitary Lie algebras constitute one of the most important classes of Lie algebras appearing in physics, as shown their application in hadron physics, the quark model or spontaneously broken symmetries. In combination with other Lie algebras, pseudo-unitary algebras are also relevant for physical phenomena, such as the quaplectic group in the Born reciprocity. Generalized Inönü–Wigner contractions of these algebras are therefore also of interest, because they allow us to relate certain model transitions, as happens for the quantum mechanical versions of the liquid drop model [1], where models exist whose dynamical group is respectively the unitary group $U(6)$ and the inhomogeneous unitary group $IU(5)$. Since the latter can be obtained by means of Inönü–Wigner contractions of $U(6)$, the corresponding models are related by a limiting process. Inhomogeneous unitary algebras have also been used to describe quantum states in relativistic theories of internal particle properties [2]. In all these problems, effective expressions for the Casimir operators of the unitary and inhomogeneous unitary algebras are needed. In contrast to the simple or reductive algebras, the computation of the invariants of inhomogeneous algebras \mathfrak{g} is not feasible using the classical theory, and various approaches

have been developed [3–6], using either the universal enveloping algebra $\mathfrak{U}(\mathfrak{g})$ explicitly or by means of analytical reductions of differential equations.

In this work we analyse the generalization of the matrix procedure of Gel'fand [7] to various inhomogeneous Lie algebras. This generalization does not involve computations in the enveloping algebras or reductions of differential equations, but focuses on the combination of the Gel'fand method for simple Lie algebras with contractions of Lie algebras [8, 9]. We first compute the Casimir operators of inhomogeneous pseudo-unitary algebras $Iu(p, q)$ by means of characteristic polynomials, and then extend the result to certain of its contractions. We also analyse the corresponding problem for inhomogeneous special pseudo-unitary algebras $I\mathfrak{su}(p, q)$, which are obtained contracting¹ either $\mathfrak{su}(p+1, q)$ or $\mathfrak{su}(p, q+1)$ [10]. However, here the contraction of the corresponding Casimir operators leads to dependence problems. To avoid this situation, we construct a special contraction $K\mathfrak{su}(p, q)$ of $\mathfrak{su}(p+1, q)$, which can be realized as an extension by a derivation of the inhomogeneous algebras $I\mathfrak{su}(p, q)$. We give a determinantal formula for the invariants of this contraction, and determine a maximal set $\{W_k\}$ of invariants which satisfy a certain constraint. Analysing how these functions W_k transform when inserted in the system of PDEs giving the invariants of $I\mathfrak{su}(p, q)$, we propose a simple analytical method to compute Casimir operators of the latter algebra. This will moreover explain why contraction of $\mathfrak{su}(p, q)$ -invariants leads to powers of the quadratic Casimir operator of $I\mathfrak{su}(p, q)$, as already observed in [10]. We finally present matrix formulae to determine the Casimir invariants of other Lie algebras, namely the inhomogeneous Lie algebras $I\mathfrak{so}(p, q)$ and two of its contractions that can be seen as a generalization of the classical Galilei and Carroll Lie algebras.

The method to obtain the invariants of a Lie algebra that we will use in this work is the analytical approach, i.e., using the representation of \mathfrak{g} by differential operators

$$\widehat{X}_i = -C_{ij}^k x_k \frac{\partial}{\partial x_j}, \quad 1 \leq i \leq n, \quad (1)$$

where $\{C_{ij}^k\}$ is the structure tensor over the basis $\{X_1, \dots, X_n\}$ and $\{x_1, \dots, x_n\}$ the dual basis. A function $F(x_1, \dots, x_n)$ is an invariant if it satisfies the system:

$$\widehat{X}_i F(x_1, \dots, x_n) = -C_{ij}^k x_k \frac{\partial}{\partial x_j} F(x_1, \dots, x_n) = 0, \quad 1 \leq i \leq n. \quad (2)$$

If F is a polynomial solution of (2), its symmetrization corresponds to a classical Casimir operator (see e.g. [3]). The number $\mathcal{N}(\mathfrak{g})$ of independent solutions of the system (non necessarily polynomials) is given by [11]:

$$\mathcal{N}(\mathfrak{g}) = \dim \mathfrak{g} - \text{rank} (C_{ij}^k x_k), \quad (3)$$

where $A(\mathfrak{g}) := (C_{ij}^k x_k)$ is the matrix which represents the commutator table over the basis $\{X_1, \dots, X_n\}$.

We also recall the elementary notions about contractions that will be needed here. Let \mathfrak{g} be a Lie algebra and $\Phi_t \in \text{Aut}(\mathfrak{g})$ a family of automorphisms of \mathfrak{g} , where $t \in \mathbb{N}$. For any $X, Y \in \mathfrak{g}$ define

$$[X, Y]_{\Phi_t} := [\Phi_t(X), \Phi_t(Y)] = \Phi_t([X, Y]). \quad (4)$$

Obviously $[X, Y]_{\Phi_t}$ are the brackets of the Lie algebra over the transformed basis. Now suppose that the limit

$$[X, Y]_{\infty} := \lim_{t \rightarrow \infty} \Phi_t^{-1} [\Phi_t(X), \Phi_t(Y)] \quad (5)$$

¹ More specifically, as a direct summand of a contraction.

exists for any $X, Y \in \mathfrak{g}$. Equation (5) defines a Lie algebra \mathfrak{g}' called the contraction of \mathfrak{g} (by Φ_t). If \mathfrak{g} and \mathfrak{g}' are nonisomorphic, then the contraction is called non-trivial. This procedure can be enlarged to contract Casimir invariants. If $F(X_1, \dots, X_n)$ is a Casimir operator of degree p of \mathfrak{g} , we can rewrite it in the new basis $\{\Phi_t(X_1), \dots, \Phi_t(X_n)\}$ and consider the limit

$$F'(X_1, \dots, X_n) := \lim_{t \rightarrow \infty} t^p F(\Phi_t(X_1), \dots, \Phi_t(X_n)). \tag{6}$$

It is straightforward to verify, using (5), that $F'(X_1, \dots, X_n)$ is a Casimir operator of the contraction. This method has become standard in the literature to obtain Casimir invariants of contractions [8, 9, 12, 13].

2. Inhomogeneous pseudo-unitary Lie algebras and their invariants

In this work we will use the basis of $u(p, q)$ given by the operators $\{E_{\mu\nu}, F_{\mu\nu}\}_{1 \leq \mu, \nu \leq p+q=n}$ with the constraints $E_{\mu\nu} + E_{\nu\mu} = 0, F_{\mu\nu} - F_{\nu\mu} = 0$. The commutation relations over this basis are

$$[E_{\mu\nu}, E_{\lambda\sigma}] = g_{\mu\lambda}E_{\nu\sigma} + g_{\mu\sigma}E_{\lambda\nu} - g_{\nu\lambda}E_{\mu\sigma} - g_{\nu\sigma}E_{\lambda\mu} \tag{7}$$

$$[E_{\mu\nu}, F_{\lambda\sigma}] = g_{\mu\lambda}F_{\nu\sigma} + g_{\mu\sigma}F_{\lambda\nu} - g_{\nu\lambda}F_{\mu\sigma} - g_{\nu\sigma}F_{\lambda\mu} \tag{8}$$

$$[F_{\mu\nu}, F_{\lambda\sigma}] = g_{\mu\lambda}E_{\nu\sigma} + g_{\nu\lambda}E_{\mu\sigma} - g_{\nu\sigma}E_{\lambda\mu} - g_{\mu\sigma}E_{\lambda\nu}, \tag{9}$$

where $(g_{\mu\nu}) = (1, \dots, 1, -1, \dots, -1)$ is a diagonal matrix. Since $u(p, q) = \mathfrak{su}(p, q) \oplus \mathbb{R}$, it follows at once that $u(p, q)$ has $(p + q)$ independent Casimir operators, one of them being $g^{\mu\mu}F_{\mu\mu}$, for corresponding to the central element, while the other $(p + q - 1)$ invariants correspond to the simple part [14–16].

Proposition 1. For $N = p + q \geq 2$, a maximal set of independent Casimir invariants of $u(p, q)$ is given by the coefficients C_k of the characteristic polynomial $|\mathfrak{i}A_{p,q} - \lambda \mathfrak{I}d_N| = \lambda^N + \sum_{k=1}^N C_k \lambda^{N-k}$, where

$$A_{p,q} = \begin{pmatrix} -\mathfrak{i}f_{11} & \dots & -g_{jj}(e_{1j} + \mathfrak{i}f_{1j}) & \dots & -g_{NN}(e_{1N} + \mathfrak{i}f_{1N}) \\ \vdots & & \vdots & & \vdots \\ e_{1j} - \mathfrak{i}f_{1j} & \dots & -\mathfrak{i}g_{jj}f_{jj} & \dots & -g_{NN}(e_{jN} + \mathfrak{i}f_{jN}) \\ \vdots & & \vdots & & \vdots \\ e_{1N} - \mathfrak{i}f_{1N} & \dots & g_{jj}(e_{jN} - \mathfrak{i}f_{jN}) & \dots & -\mathfrak{i}g_{NN}f_{NN} \end{pmatrix} \tag{10}$$

and $\mathfrak{i} = \sqrt{-1}$. Moreover, $\deg(C_k) = k$ for $k = 1 \dots p + q = N$.

In particular, if $g_{ii} = 1$ for all i , we obtain the unitary algebra $u(n)$. Using the notation $a^{ij} = (e_{ij} + \mathfrak{i}f_{ij})$ with the constraint $a^{ij} = -a^{ji*}$, it was shown in [18, chapter 7, p 239] that the Casimir operators of $u(n)$ are obtained from the secular equation

$$\det \|a^{ij}J_{ij} - \lambda I_n\| = (-\lambda)^{n-r} \phi_r(a^{ij}), \tag{11}$$

where

$$\phi_r(a^{ij}) = \frac{1}{r!(n-r)!} \varepsilon_{i_1 \dots i_n} a^{i_1 j_1} \dots a^{i_r j_r} \varepsilon_{j_1 \dots j_r}. \tag{12}$$

Here J_{ij} is the fundamental matrix having entry 1 at position (i, j) and zero elsewhere, and ε the skew Levi-Civita tensor. Formula (10) is nothing but the corresponding adaptation to the different real forms of $u(N)$ over the basis (7)–(9).

There is an interesting consequence of this formulation, namely, that the determinant of $A_{p,q}$ satisfies the following identity:

$$\det(A_{p,q}) = \prod_{j=2}^N g_{jj} \det(A_{N,0}). \tag{13}$$

This implies that all pseudo-unitary algebras $u(p, q)$ have the Casimir invariant of order N in common, while those of lower degree depend essentially on the index and signature of the matrix $(g_{\mu\mu})$.

As shown first by Rosen, inhomogeneous unitary Lie algebras $Iu(p, q)$ can be obtained from an Inönü–Wigner contraction of the unitary algebras $u(p + 1, q)$ or $u(p, q + 1)$. We will use also this contraction to obtain the matrix formulae giving the Casimir operators of $Iu(p, q)$, but without working explicitly with elements of the enveloping algebra. Our approach uses only matrices, and the invariants will follow from characteristic polynomials or combinations of them. To this extent, we must first fix some notation. We consider the ordered basis F_{ii}, E_{ij}, F_{ij} for $1 \leq i, j \leq N$. Without loss of generality, we can suppose that for the matrix (10) we have $g_{11} = 1$. We consider the automorphism Φ of $u(p, q)$ given by

$$\begin{aligned} \Phi(F_{11}) &= \frac{1}{t} F_{11}, & \Phi(F_{\mu\mu}) &= F_{\mu\mu} (\mu \geq 2); & \Phi(E_{1\sigma}) &= \frac{1}{t} E_{1\sigma} (\sigma \geq 2) \\ \Phi(F_{1\sigma}) &= \frac{1}{t} F_{1\sigma} (\sigma \geq 2) & \Phi(E_{\mu\nu}) &= E_{\mu\nu}, & \Phi(F_{\mu\nu}) &= F_{\mu\nu}, \mu, \nu \neq 1. \end{aligned} \tag{14}$$

For $\mu, \nu \neq 1$ the generators $\Phi(E_{\mu\nu})$ and $\Phi(F_{\mu\nu})$ obviously generate the unitary Lie algebra $u(p - 1, q)$. The limit $\lim_{t \rightarrow \infty} \Phi^{-1}[\Phi(X), \Phi(Y)]$ exists for any pair of elements $X, Y \in u(p, q)$, and clearly defines a contraction. If we write $\Phi(F_{1\rho}) := Q_\rho$ and $\Phi(E_{1\rho}) := R_\rho$, then we have the brackets

$$\begin{aligned} [E_{\mu\nu}, R_\rho] &= g_{\mu\rho} R_\nu - g_{\nu\rho} R_\mu, & [E_{\mu\nu}, Q_\rho] &= g_{\mu\rho} Q_\nu - g_{\nu\rho} Q_\mu, \\ [F_{\mu\nu}, Q_\rho] &= -g_{\nu\rho} R_\mu - g_{\mu\rho} R_\nu, & [F_{\mu\nu}, R_\rho] &= g_{\mu\rho} Q_\nu + g_{\nu\rho} Q_\mu, \end{aligned} \tag{15}$$

and this coincides exactly with the standard representation of $u(p - 1, q)$ as given in [10]. Observe that in particular $\Phi(F_{11})$ commutes with any other generator. Therefore the contraction is isomorphic to the direct sum $Iu(p - 1, q) \oplus \mathbb{R}$.

Let $N + 1 = p + q + 1 \geq 2$ and consider the contraction $u(p, q) \rightsquigarrow Iu(p - 1, q) \oplus \langle F_{11} \rangle$ determined by (14). We next give with an analogous formula for the inhomogeneous Lie algebra $Iu(p - 1, q)$.

Proposition 2. *In the preceding conditions, a maximal set of independent Casimir operators D_k of $Iu(p - 1, q)$ is obtained from the determinant $\Delta(\lambda) = \sum_{k=2}^N D_k \lambda^{N-k}$ defined by*

$$\lim_{t \rightarrow \infty} \frac{1}{t^2} \begin{vmatrix} f_{11}t - \lambda & \dots & -g_{jj}(ie_{1j} - f_{1j})t & \dots & -g_{N+1,N+1}(ie_{1,N+1} - f_{1,N+1})t \\ \vdots & & \vdots & & \vdots \\ (ie_{1j} + f_{1j})t & \dots & g_{jj}f_{jj} - \lambda & \dots & -g_{N+1,N+1}(ie_{j,N+1} - f_{j,N+1}) \\ \vdots & & \vdots & & \vdots \\ (ie_{1,N+1} + f_{1,N+1})t & \dots & g_{jj}(ie_{j,N+1} + f_{j,N+1}) & \dots & g_{N+1,N+1}f_{N+1,N+1} - \lambda \end{vmatrix}. \tag{16}$$

The preceding formula is obtained starting from (10) and implementing the contraction to the algebra $Iu(p, q) \oplus \mathbb{R}$, scaling the determinant to avoid divergencies (see equation (6)).

The independence of the contracted invariants is proved directly. It follows from equation (10) that for any $2 \leq j \leq N$ the following identities hold:

$$\frac{\partial^3 C_j}{\partial e_{1\mu} \partial e_{1\nu} \partial e_{1\lambda}} = \frac{\partial^3 C_j}{\partial e_{1\mu} \partial e_{1\nu} \partial f_{1\lambda}} = \frac{\partial^3 C_j}{\partial f_{1\mu} \partial f_{1\nu} \partial e_{1\lambda}} = \frac{\partial^3 C_j}{\partial f_{1\mu} \partial f_{1\nu} \partial f_{1\lambda}} = 0. \tag{17}$$

That means that any invariant² of $u(p, q)$ can be written as a quadratic polynomial in the variables $\{e_{1\mu}, f_{1\mu}\}$ with coefficients in the polynomial ring $\mathbb{R}[e_{\mu\nu}, f_{\mu\nu}]_{2 \leq \mu, \nu \leq N}$. Moreover, it is straightforward to verify that no function depending only on the variables $\{e_{\mu\nu}, f_{\mu\nu}\}$ with $\mu, \nu \neq 1$ is an invariant of $Iu(p, q)$, because of the action of the reductive part $u(p - 1, q)$ over the representation given in (15). From these facts it follows that the functions D_k are algebraically independent homogeneous polynomials, and therefore their symmetrization are independent Casimir invariants [17]. By algebraic manipulation, we can reduce and simplify the determinant (16) in order to obtain an expression that no longer involves the contraction parameter t . This is done by applying the usual determinantal methods, like the Laplace expansion. After some decompositions and computations, we obtain that

$$\Delta = |B_{p,q} - \lambda \text{Id}_N| + \lambda |(B_{p,q})_{11} - \lambda \text{Id}_{N-1}|, \tag{18}$$

where

$$B_{p,q} := \begin{pmatrix} 0 & \cdots & g_{jj}(f_{1j} - ie_{1j}) & \cdots & g_{N+1,N+1}(f_{1,N+1} - ie_{1,N+1}) \\ \vdots & & \vdots & & \vdots \\ (ie_{1j} + f_{1j}) & \cdots & g_{jj}f_{jj} & \cdots & g_{N+1,N+1}(f_{j,N+1} - ie_{j,N+1}) \\ \vdots & & \vdots & & \vdots \\ (ie_{1,N+1} + f_{1,N+1}) & \cdots & g_{jj}(ie_{j,N+1} + f_{j,N+1}) & \cdots & g_{N+1,N+1}f_{N+1,N+1} \end{pmatrix}, \tag{19}$$

and $(B_{p,q})_{11}$ is the minor obtained by deleting the first row and column. Now, replacing e_{1j} by r_j and f_{1j} by q_j , equation (19) gives the Casimir operators of $Iu(p - 1, q)$ over the usual basis $\{E_{\mu\nu}, F_{\mu\nu}, Q_\rho, R_\rho\}$. The advantage of formula (19) with respect to equation (16) is that the invariants are expressed in the usual basis of $Iu(p, q)$, without taking into account that this algebra arises as a contraction of the unitary algebra $u(p + 1, q)$, and therefore avoiding the computation of limits.

3. A comment concerning the special pseudo-unitary algebras

Formula (10) can be easily adapted to the special algebras. For convenience we take the Cartan subalgebra spanned by the vectors $H_\mu = g_{\mu+1,\mu+1}F_{\mu\mu} - g_{\mu\mu}F_{\mu+1,\mu+1}$ for $\mu = 1 \cdots p + q - 1$. The centre of $u(p, q)$ is obviously generated by $g^{\mu\mu}F_{\mu\mu}$, which coincides with the first order invariant obtained previously.

Proposition 3. *For $N \geq 2$, a maximal set of independent Casimir invariants of $su(p, q)$ (where $p + q = N$) is given by the coefficients D_k of the characteristic polynomial*

² Of degree $d \geq 2$.

$|iA_{p,q} - \lambda I d_N| = \lambda^N + \sum_{k=2}^N D_k \lambda^{N-k}$, where

$$A_{p,q} = \begin{pmatrix} -iY_1 & \cdots & -g_{\mu\mu}(e_{1\mu} + if_{1\mu}) & \cdots & -g_{NN}(e_{1N} + if_{1N}) \\ \vdots & & \vdots & & \vdots \\ e_{1\mu} - if_{1\mu} & \cdots & -iY_\mu & \cdots & -g_{NN}(e_{\mu N} + if_{\mu N}) \\ \vdots & & \vdots & & \vdots \\ e_{1N} - i_{1N} & \cdots & g_{\mu\mu}(e_{\mu N} - if_{\mu N}) & \cdots & -iY_N \end{pmatrix}, \quad (20)$$

and

$$Y_\mu = \sum_{v=1}^{\mu-1} \frac{-v}{N} g_{vv} g_{v+1,v+1} h_v + \sum_{v=\mu}^{N-1} \frac{N-v}{N} g_{vv} g_{v+1,v+1} h_v, \quad 1 \leq \mu \leq N. \quad (21)$$

Moreover, $\deg(D_k) = k$ for $k = 2 \cdots p + q = N$.

The interest of the preceding formula (20) is its application to obtain closed expressions for the Casimir operators of some contractions of $\mathfrak{su}(p, q)$. Consider for instance the automorphism

$$\begin{aligned} \Phi(H_\mu) &= H_\mu (\mu \geq 1), & \Phi(E_{1\sigma}) &= \frac{1}{t} E_{1\sigma} (\sigma \geq 2), & \Phi(F_{1\sigma}) &= \frac{1}{t} F_{1\sigma} (\sigma \geq 2), \\ \Phi(E_{\mu\nu}) &= E_{\mu\nu}, & \Phi(F_{\mu\nu}) &= F_{\mu\nu}, \mu, \nu \neq 1. \end{aligned} \quad (22)$$

For any pair of elements X, Y of the algebra $\lim_{t \rightarrow \infty} \Phi^{-1}[\Phi(X), \Phi(Y)]$ exists, and defines a Lie algebra which will be denoted by $K\mathfrak{su}(p-1, q)$.³ It follows at once from (22) that $K\mathfrak{su}(p-1, q)$ is isomorphic to an extension of degree one of the inhomogeneous algebra $I\mathfrak{su}(p-1, q)$. In fact, observe that the action of the Cartan subalgebra \mathfrak{h} over the $E_{\mu,v}$ and $F_{\mu,v}$ is not changed by the contraction. If we consider the element $H' \in \mathfrak{h}$ defined by

$$H' = \sum_{\mu=1}^{p+q-1} (p+q-\mu) \prod_{\rho \neq \mu, \mu+1} g_{\rho\rho} H_\mu, \quad (23)$$

it is not difficult to see that taking the new basis of \mathfrak{h} generated by $\{H', H_2, \dots, H_{p+q-1}\}$ we have $[H', E_{\mu\nu}] = [H', F_{\mu\nu}] = 0$ and

$$[H', E_{1\mu}] = -2(p+q) \prod_{\rho=1}^{p+q} g_{\rho\rho} F_{1\mu}, \quad [H', F_{1\mu}] = 2(p+q) \prod_{\rho=1}^{p+q} g_{\rho\rho} E_{1\mu} \quad (24)$$

for $2 \leq \mu < v$. This shows that $[H', \mathfrak{su}(p-1, q)] = 0$. Using this fact, it is straightforward to verify that the linear mapping $\varphi : I\mathfrak{su}(p-1, q) \rightarrow I\mathfrak{su}(p-1, q)$ defined by $\varphi(X) = [H', X]$ is a derivation of the inhomogeneous algebra $I\mathfrak{su}(p-1, q)$. This result allows us to determine the Casimir operators of the extension using the preceding formulae. The proof is completely analogous to that of proposition 2.

Proposition 4. For $N = p + q \geq 2$, a maximal set of independent Casimir invariants of $K\mathfrak{su}(p-1, q)$ over the basis $\{H', H_2, \dots, H_{N-1}, E_{\mu\nu}, F_{\mu\nu}\}$ is given by the coefficients W_k of the polynomial

$$|iC_{p,q} - \lambda I d_N| + \lambda |i(C_{p,q})_{11} - \lambda I d_{N-1}| = \sum_{k=2}^N W_k \lambda^{N-k}, \quad (25)$$

³ The Levi part, which is isomorphic to $\mathfrak{su}(p-1, q)$, is generated by $\{H_\mu, E_{\mu\nu}, F_{\mu\nu}\}_{2 \leq \mu < \nu}$.

where

$$C_{p,q} = \begin{pmatrix} 0 & \cdots & -g_{\mu\mu}(e_{1\mu} + if_{1\mu}) & \cdots & -g_{NN}(e_{1N} + if_{1N}) \\ \vdots & & \vdots & & \vdots \\ e_{1\mu} - if_{1\mu} & \cdots & -iY_\mu & \cdots & -g_{NN}(e_{\mu N} + if_{\mu N}) \\ \vdots & & \vdots & & \vdots \\ e_{1N} - if_{1N} & \cdots & g_{\mu\mu}(e_{\mu N} - if_{\mu N}) & \cdots & -iY_N \end{pmatrix}, \tag{26}$$

and

$$Y_\mu = - \prod_{\rho=1}^N g_{\rho\rho} \frac{h'}{N(N-1)} + \sum_{v=2}^{\mu-1} \frac{1-v}{N-1} g_{vv} g_{v+1,v+1} h_v + \sum_{v=\mu}^{N-1} \frac{N-v}{N-1} g_{vv} g_{v+1,v+1} h_v \tag{27}$$

for $2 \leq \mu \leq N$. Moreover, $\deg(W_k) = k$ for $k = 2 \dots p + q = N$.

The functions W_k have a structural property that will be of interest in the analysis of the Casimir operators of inhomogeneous algebra $I\mathfrak{su}(p, q)$.

Proposition 5. For any $N \geq 2$, the functions W_k of (25) satisfy the equation

$$N(N-1) \frac{\partial W_k}{\partial h'} - (N+1-k) \prod_{\rho=1}^N g_{\rho\rho} W_{k-1} = 0. \tag{28}$$

Proof. By (25), the W_k are the result of a combination of the characteristic polynomial of the matrix $iC_{p,q}$ and the minor $i(C_{p,q})$ obtained deleting the first row and column. Therefore, using the expansion properties of determinants, it follows easily that any W_k is obtained from the sum of the minors of order k of $iC_{p,q}$ that contain the first column and row, i.e.,

$$W_k = \sum_{2 < \mu_1 < \dots < \mu_{k-1}} M(1, \mu_1, \dots, \mu_{k-1}), \tag{29}$$

where $M(1, \mu_1, \dots, \mu_{k-1})$ is the matrix given by

$$\begin{pmatrix} 0 & -g_{\mu_1\mu_1}(ie_{1\mu_1} - f_{1\mu_1}) & \cdots & -g_{\mu_{k-1},\mu_{k-1}}(ie_{1\mu_{k-1}} - f_{1\mu_{k-1}}) \\ (ie_{1\mu_1} + f_{1\mu_1}) & Y_{\mu_1} & \cdots & -g_{\mu_{k-1},\mu_{k-1}}(ie_{\mu_1\mu_{k-1}} - f_{\mu_1\mu_{k-1}}) \\ \vdots & \vdots & & \vdots \\ (ie_{1\mu_{k-1}} + f_{1\mu_{k-1}}) & g_{\mu_1\mu_1}(ie_{\mu_1\mu_{k-1}} + f_{\mu_1\mu_{k-1}}) & \cdots & Y_{\mu_{k-1}} \end{pmatrix}. \tag{30}$$

Now the variable h' is contained only in the elements Y_μ of the diagonal, which implies that

$$\begin{aligned} \frac{\partial W_k}{\partial h'} &= - \sum_{2 < \mu_1 < \dots < \mu_{k-1}} \frac{\partial M(1, \mu_1, \dots, \mu_{k-1})}{\partial h'} = - \sum_{2 < \mu_1 < \dots < \mu_{k-1}} \sum_{j=\mu_1}^{\mu_{k-1}} \frac{\partial Y_j}{\partial h'} M(1, \mu_1, \dots, \mu_{k-2}) \\ &= \frac{(N+1-k)}{N(N-1)} \prod_{\rho=1}^N g_{\rho\rho} \sum_{2 < \mu_1 < \dots < \mu_{k-2}} M(1, \mu_1, \dots, \mu_{k-2}) = \frac{(N+1-k)}{N(N-1)} \prod_{\rho=1}^N g_{\rho\rho} W_{k-1}. \end{aligned} \tag{31}$$

□

Since $K\mathfrak{su}(p-1, q)$ is a contraction of $\mathfrak{su}(p, q)$, by transitivity of contractions it follows that $K\mathfrak{su}(p-1, q)$ also contracts onto $I\mathfrak{su}(p-1, q) \oplus \langle H' \rangle$, and we obtain the chain of contractions

$$\mathfrak{su}(p, q) \rightsquigarrow K\mathfrak{su}(p-1, q) \rightsquigarrow I\mathfrak{su}(p-1, q) \oplus \langle H' \rangle. \tag{32}$$

The question is whether the procedure developed to compute the Casimir invariants can be adapted to the special inhomogeneous case. The following example shows that in general, the contraction of the invariants forming a fundamental set of invariants of $\mathfrak{su}(p, q)$ (or $K\mathfrak{su}(p - 1, q)$) does not generate a complete set of invariants of the contraction. Consider the Lie algebra $\mathfrak{su}(3, 1)$. This algebra has three invariants C_2, C_3, C_4 which are obtained from the characteristic polynomial of the matrix

$$D = \begin{pmatrix} \frac{3}{4}h_1 + \frac{1}{2}h_2 - \frac{1}{4}h_3 & -ie_{12} + f_{12} & -ie_{13} + f_{13} & ie_{14} - f_{14} \\ ie_{12} + f_{12} & -\frac{1}{4}h_1 + \frac{1}{2}h_2 - \frac{1}{4}h_3 & -ie_{23} + f_{23} & ie_{24} - f_{24} \\ ie_{13} + f_{13} & ie_{23} + f_{23} & -\frac{1}{4}h_1 - \frac{1}{2}h_2 - \frac{1}{4}h_3 & ie_{34} - f_{34} \\ ie_{14} + f_{14} & ie_{24} + f_{24} & ie_{34} + f_{34} & -\frac{1}{4}h_1 - \frac{1}{2}h_2 + \frac{3}{4}h_3 \end{pmatrix}. \tag{33}$$

We obtain the inhomogeneous algebra $I\mathfrak{su}(2, 1)$ from the contraction determined by the automorphism

$$\begin{aligned} \Phi(H_1) &= \frac{1}{t}H_1, & \Phi(F_\mu) &= F_\mu (\mu \geq 2), & \Phi(E_{1\sigma}) &= \frac{1}{t}E_{1\sigma} (\sigma \geq 2), \\ \Phi(F_{1\sigma}) &= \frac{1}{t}F_{1\sigma} (\sigma \geq 2), & \Phi(E_{\mu\nu}) &= E_{\mu\nu}, & \Phi(F_{\mu\nu}) &= F_{\mu\nu}, \mu, \nu \neq 1. \end{aligned} \tag{34}$$

However, since we are rescaling also a generator of the Cartan subalgebra, it will follow that the contraction of the Casimir operators $C_i(h_\mu, e_{\mu\nu}, f_{\mu\nu})$ for $i = 2, 3, 4$ are functions of the quadratic operator $I_2 = g_{\mu\mu}(e_{1\mu}^2 + f_{1\mu}^2)$ of $I\mathfrak{su}(2, 1)$ and h_1 .⁴ More specifically, for $k = 2, 3, 4$ we have

$$\lim_{t \rightarrow \infty} \frac{1}{t^k} C_k = -\alpha_k I_2 h_1^{k-2}, \tag{35}$$

where $\alpha_2 = 1, \alpha_3 = \frac{1}{2}, \alpha_4 = \frac{3}{16}$. Therefore the characteristic polynomial of the matrix (33) does not provide a maximal set of invariants of $I\mathfrak{su}(2, 1)$. However, it is easy to obtain an independent invariant using (35). This third invariant, which must depend on all variables of $I\mathfrak{su}(2, 1)$ [6], can be obtained for example contracting the invariant $C_2 C_4 - \frac{1}{4} C_3^2$, and provides a sixth order Casimir operator of $I\mathfrak{su}(2, 1)$ which is not a power of I_2 . A similar result holds if we start from the invariants of $K\mathfrak{su}(2, 1)$. We can however proceed differently. Instead of contracting the invariants W_k , we evaluate them in the system (2) corresponding to $I\mathfrak{su}(2, 1)$ over the basis $\{H_2, H_3, E_{\mu\nu}, F_{\mu\nu}\}$. We obtain

$$\begin{aligned} \widehat{H}_\mu(W_i) &= 0 (\mu = 2, 3), & \widehat{E}_{\mu\nu}(W_2) &= 0, & \widehat{F}_{\mu\nu}(W_2) &= 0, \\ \widehat{E}_{\mu\nu}(W_3) &= -\frac{4}{3} f_{\mu\nu} W_2, & \widehat{F}_{\mu\nu}(W_3) &= \frac{4}{3} e_{\mu\nu} W_2, & \widehat{E}_{\mu\nu}(W_4) &= -\frac{2}{3} f_{\mu\nu} W_3. \end{aligned} \tag{36}$$

Now let $u = W_3$ and $v = W_4$. For any $X \in \{E_{\mu\nu}, F_{\mu\nu}\}$ we have

$$\frac{\widehat{X}(v)}{\widehat{X}(u)} = \frac{W_3}{2W_2} = \frac{u}{2W_2}.$$

If we now consider the differential equation

$$\frac{\partial F}{\partial u} + \frac{u}{2W_2} \frac{\partial F}{\partial v} = 0, \tag{37}$$

it follows easily that the solutions are generated by the function $(W_2 v - \frac{1}{4} u^2) W_2^{-1} = (W_2 W_4 - \frac{1}{4} W_3^2) W_2^{-1}$. Since W_2 is already an invariant of $I\mathfrak{su}(2, 1)$, this implies that

⁴ Observe that in the contraction this becomes a central element.

$I = 4W_2W_4 - W_3^2$ is an invariant of the algebra of degree six. A short computation shows that it coincides with $\lim_{t \rightarrow \infty} \frac{1}{t^4} (C_2C_4 - \frac{1}{4}C_3^2)$. This fact is not casual and can be formulated in any dimension.

Theorem 1. *Let $p + q = N \geq 2$ and W_k be the invariants of the Lie algebra $K\mathfrak{su}(p - 1, q)$. A function $G(W_2, \dots, W_N)$ is an invariant of the inhomogeneous Lie algebra $I\mathfrak{su}(p - 1, q)$ if and only if*

$$\frac{\partial G(W_2, \dots, W_N)}{\partial h'} = 0.$$

Proof. Let us consider the bases of $K\mathfrak{su}(p - 1, q)$ and $I\mathfrak{su}(p - 1, q)$ used previously, and let $\{\widehat{H}', \widehat{H}_\mu, \widehat{E}_{\mu\nu}, \widehat{F}_{\mu\nu}\}$ and $\{\widetilde{H}', \widetilde{H}_\mu, \widetilde{E}_{\mu\nu}, \widetilde{F}_{\mu\nu}\}$ be the representation by differential operators of $K\mathfrak{su}(p - 1, q)$, respectively $I\mathfrak{su}(p - 1, q) \oplus \langle H' \rangle$. Since $K\mathfrak{su}(p - 1, q)$ is an extension by a derivation of $I\mathfrak{su}(p - 1, q)$ acting trivially on the Levi part, for the invariants W_k we obtain the relations

$$\widehat{H}_\mu(W_\lambda) = \widetilde{H}_\mu(W_\lambda) = 0, \quad 1 \leq \mu \leq N - 1 \tag{38}$$

$$\widehat{E}_{\mu\nu}(W_\lambda) = \widetilde{E}_{\mu\nu}(W_\lambda) = 0, \quad 2 \leq \mu, \quad \nu \leq N \tag{39}$$

$$\widehat{F}_{\mu\nu}(W_\lambda) = \widetilde{F}_{\mu\nu}(W_\lambda) = 0, \quad 2 \leq \mu, \quad \nu \leq N \tag{40}$$

$$\widehat{E}_{1\mu}(W_\lambda) = \widetilde{E}_{1\mu}(W_\lambda) + 2N \prod_{\rho=1}^N g_{\rho\rho} \frac{\partial W_\lambda}{\partial h'} f_{1\mu} = 0, \quad 2 \leq \mu \leq N \tag{41}$$

$$\widehat{F}_{1\mu}(W_\lambda) = \widetilde{F}_{1\mu}(W_\lambda) - 2N \prod_{\rho=1}^N g_{\rho\rho} \frac{\partial W_\lambda}{\partial h'} e_{1\mu} = 0, \quad 2 \leq \mu \leq N. \tag{42}$$

For W_2 we obviously have $\frac{\partial W_2}{\partial h'} = 0$, we thus recover the quadratic invariant. For any $\lambda \geq 3$ the previous system, combined with proposition 5, provides the following expressions for any $2 \leq \mu \leq N$:

$$\widetilde{E}_{1\mu}(W_\lambda) = -2N \prod_{\rho=1}^N g_{\rho\rho} \frac{\partial W_\lambda}{\partial h'} = -\frac{2(N + 1 - k)}{(N - 1)} W_{\lambda-1} f_{1\mu}, \tag{43}$$

$$\widetilde{F}_{1\mu}(W_\lambda) = 2N \prod_{\rho=1}^N g_{\rho\rho} \frac{\partial W_\lambda}{\partial h'} = \frac{2(N + 1 - k)}{(N - 1)} W_{\lambda-1} e_{1\mu}. \tag{44}$$

As a consequence of these identities, an invariant W of $K\mathfrak{su}(p - 1, q)$ is also an invariant of the inhomogeneous algebra $I\mathfrak{su}(p - 1, q)$ if and only if $\frac{\partial W}{\partial h'} = 0$. \square

Observe that equations (43) and (44) can be used to compute additional invariants of $I\mathfrak{su}(p - 1, q)$ using differential equations of the type (37). The fact that the resulting function must be independent of h' will impose some restrictions on the degree.

4. On the applicability of the method

Although the method to compute Casimir invariants by means of characteristic polynomials has been developed for the Lie algebras $u(p, q)$ and their contractions, the argument is not

exclusive for the unitary case, and does not depend on the special shape of their invariants. In this section we give matrix formulae for the invariants of other inhomogeneous Lie algebras of physical importance. More specifically, we apply the method to the inhomogeneous pseudo-orthogonal Lie algebras, a double inhomogeneous algebra and a semidirect product of a Heisenberg and an orthogonal Lie algebra, which can both be obtained from contractions of $I\mathfrak{so}(p, q)$ and were studied in [4].

4.1. Inhomogeneous pseudo-orthogonal Lie algebras

The inhomogeneous pseudo-orthogonal algebra $I\mathfrak{so}(p, q)$ with $N = p + q$ is given by the $\frac{1}{2}N(N + 1)$ operators $E_{\mu\nu} = -E_{\nu\mu}, P_\mu$, satisfying

$$[E_{\mu\nu}, E_{\lambda\sigma}] = g_{\mu\lambda}E_{\nu\sigma} + g_{\mu\sigma}E_{\lambda\nu} - g_{\nu\lambda}E_{\mu\sigma} - g_{\nu\sigma}E_{\lambda\mu}$$

$$[E_{\mu\nu}, P_\rho] = g_{\mu\rho}P_\nu - g_{\nu\rho}P_\mu,$$

where $g = \text{diag}(1, \dots, 1, -1, \dots, -1)$. Moreover, in [3] it was shown that $\mathcal{N}(I\mathfrak{so}(p, q)) = \lceil \frac{p+q+1}{2} \rceil$.

Proposition 6. A maximal set of Casimir operators of $I\mathfrak{so}(p, q)$ is given by the coefficients C_k of $P(T)$ defined by

$$P(T) := |B_{p,q} - T\text{Id}_{N+1}| + T|(B_{p,q})_{(N+1,N+1)} - T\text{Id}_N|, \tag{45}$$

where

$$B_{p,q} := \begin{pmatrix} 0 & \cdots & -g_{jj}e_{1j} & \cdots & -g_{NN}e_{1,N} & p_1T \\ \vdots & & \vdots & & \vdots & \\ e_{1j} & \cdots & 0 & \cdots & -g_{NN}e_{j,N} & p_jT \\ \vdots & & \vdots & & \vdots & \\ e_{1,N} & \cdots & g_{jj}e_{j,N} & \cdots & 0 & p_NT \\ -p_1 & & -g_{jj}P_j & & -g_{NN}P_N & 0 \end{pmatrix}. \tag{46}$$

The proof is quite similar to that of proposition 2. Here we use the contraction of $\mathfrak{so}(p + 1, q)$ over $I\mathfrak{so}(p, q)$, as done in [10]. In contrast to the unitary case, where the matrices involved are complex, here no problem of dependence in the contracted invariants is observed.

4.2. The double inhomogeneous algebra $AG(p, q)$

The Lie algebra $AG(p, q)$ with $N = p + q + 1$ is given by the $\frac{1}{2}(N + 1)N$ operators $E_{\mu\nu} = -E_{\nu\mu}, P_\mu, G_\nu, P_0$, satisfying the nontrivial commutation relations:

$$[E_{\mu\nu}, E_{\lambda\sigma}] = g_{\mu\lambda}E_{\nu\sigma} + g_{\mu\sigma}E_{\lambda\nu} - g_{\nu\lambda}E_{\mu\sigma} - g_{\nu\sigma}E_{\lambda\mu}, \quad [P_0, Q_\rho] = -g_{N-1,N-1}P_\rho, \tag{47}$$

$$[E_{\mu\nu}, P_\rho] = g_{\mu\rho}P_\nu - g_{\nu\rho}P_\mu, \quad [E_{\mu\nu}, Q_\rho] = g_{\mu\rho}Q_\nu - g_{\nu\rho}Q_\mu,$$

where $1 \leq \mu, \nu, \rho, \leq N - 1$ and $(g_{\mu\nu})$ is the preceding diagonal matrix. It can be easily shown that this algebra is obtained by an Inönü–Wigner contraction of the inhomogeneous Lie algebra $I\mathfrak{so}(p + 1, q)$, considering the automorphism

$$\Psi(E_{\mu\nu}) = E_{\mu\nu}(1 \leq \mu, \nu \leq N - 1), \quad \Psi(E_{\mu N}) = \frac{1}{t}E_{\mu N} := G_\mu, \tag{48}$$

$$\Psi(P_\mu) = \frac{1}{t}P_\mu(1 \leq \mu \leq N - 1), \quad \Psi(P_N) = P_N := P_0.$$

This contraction preserves the number of invariants, i.e., $\mathcal{N}(AG(p, q)) = \lceil \frac{N+1}{2} \rceil$. Moreover, it should be remarked that $AG(p, q)$ is a double inhomogeneous Lie algebra with semidirect product structure

$$AG(p, q) = II\mathfrak{so}(p, q) = (\mathfrak{so}(p, q) \overrightarrow{\oplus} I_{N-1}) \overrightarrow{\oplus} I_N, \tag{49}$$

I_N being the standard vector representation. In particular, for $q = 0$ we obtain the proper Galilei algebra $II\mathfrak{so}(N - 1)$.

Proposition 7. *A maximal set of Casimir operators of $AG(p, q)$ is given by the coefficients C_k of the determinant $P(T)$ defined by*

$$P(T) = \begin{vmatrix} -T & \cdots & -g_{\mu\mu}e_{1\mu} & \cdots & -g_{N-1,N-1}e_{1,N-1} & -g_{N,N}q_1 & p_1T \\ \vdots & & \vdots & & \vdots & & \vdots \\ e_{1\mu} & \cdots & -T & \cdots & -g_{N-1,N-1}e_{\mu,N-1} & -g_{N,N}q_\mu & p_\mu T \\ \vdots & & \vdots & & \vdots & & \vdots \\ e_{1,N-1} & \cdots & g_{\mu\mu}e_{\mu,N-1} & \cdots & -T & -g_{N,N}q_{N-1} & p_{N-1}T \\ q_1 & & g_{\mu\mu}q_\mu & & g_{N-1,N-1}q_{N-1} & 0 & 0 \\ -p_1 & & -g_{\mu\mu}p_\mu & & -g_{N-1,N-1}p_{N-1} & 0 & 0 \end{vmatrix} + \begin{vmatrix} -T & \cdots & 0 & \cdots & 0 & p_1T \\ \vdots & & \vdots & & \vdots & \vdots \\ 0 & \cdots & -T & \cdots & 0 & p_\mu T \\ \vdots & & \vdots & & \vdots & \vdots \\ 0 & \cdots & 0 & \cdots & -T & p_{N-1}T \\ -p_1 & \cdots & -g_{\mu\mu}p_\mu & \cdots & -g_{N-1,N-1}p_{N-1} & 0 \end{vmatrix}. \tag{50}$$

4.3. The Lie algebra $C(p, q)$

The Lie algebra $C(p, q)$ with $N = p + q + 1$ is also obtained from a contraction of $I\mathfrak{so}(p + 1, q)$ by means of the automorphism defined by

$$\begin{aligned} \Psi(E_{\mu\nu}) &= E_{\mu\nu} (1 \leq \mu, \nu \leq N - 1), & \Psi(E_{\mu N}) &= \frac{1}{t} E_{\mu N} := G_\mu, \\ \Psi(P_\mu) &= \frac{1}{t} P_\mu (1 \leq \mu \leq N - 1), & \Psi(P_N) &= \frac{1}{t^2} P_N := P_0. \end{aligned} \tag{51}$$

Over the basis $\{E_{\mu\nu}, P_\mu, G_\nu, P_0\}$, the nontrivial commutation relations are

$$\begin{aligned} [E_{\mu\nu}, E_{\lambda\sigma}] &= g_{\mu\lambda}E_{\nu\sigma} + g_{\mu\sigma}E_{\lambda\nu} - g_{\nu\lambda}E_{\mu\sigma} - g_{\nu\sigma}E_{\lambda\mu}, & [P_\rho, Q_\rho] &= -g_{\rho\rho}P_0, \\ [E_{\mu\nu}, P_\rho] &= g_{\mu\rho}P_\nu - g_{\nu\rho}P_\mu, & [E_{\mu\nu}, Q_\rho] &= g_{\mu\rho}Q_\nu - g_{\nu\rho}Q_\mu, \end{aligned} \tag{52}$$

where $1 \leq \mu, \nu, \rho, \leq N - 1$. This algebra is isomorphic to the semidirect product of $\mathfrak{so}(p, q)$ and the Heisenberg Lie algebra \mathfrak{h}_N , and therefore can be written as

$$C(p, q) = II'\mathfrak{so}(p, q) = (\mathfrak{so}(p, q) \overrightarrow{\oplus} I'_{N-1}) \overrightarrow{\oplus} I_N, \tag{53}$$

where the orthogonal algebra acts on I_N through the vector representation and on I'_{N-1} through the contragredient of the vector representation. For $q = 0$ we obtain the generalization of the proper Carroll algebra $C(p, 0) = II'\mathfrak{so}(N - 1)$. In particular it follows from (53) that $\mathcal{N}(C(p, q)) = \lceil \frac{N+1}{2} \rceil$.

Proposition 8. A maximal set of Casimir operators of $C(p, q)$ is given by the coefficients C_k of the determinant $P(T)$ defined by

$$P(T) = \begin{vmatrix} -T & \cdots & -g_{\mu\mu}e_{1\mu} & \cdots & -g_{N-1,N-1}e_{1,N-1} & -g_{N,N}q_1 & p_1T \\ \vdots & & \vdots & & \vdots & & \\ e_{1\mu} & \cdots & -T & \cdots & -g_{N-1,N-1}e_{\mu,N-1} & -g_{N,N}q_\mu & p_\mu T \\ \vdots & & \vdots & & \vdots & & \\ e_{1,N-1} & \cdots & g_{\mu\mu}e_{\mu,N-1} & \cdots & -T & -g_{N,N}q_{N-1} & p_{N-1}T \\ q_1 & & g_{\mu\mu}q_\mu & & g_{N-1,N-1}q_{N-1} & 0 & p_0T \\ -p_1 & & -g_{\mu\mu}p_\mu & & -g_{N-1,N-1}p_{N-1} & -g_{N,N}p_0 & 0 \end{vmatrix}. \quad (54)$$

5. Concluding remarks

We have seen that, in analogy with the classical algebras, the Casimir operators of inhomogeneous pseudo-unitary Lie algebras $Iu(p, q)$ can be computed using only characteristic polynomials, avoiding the use of the enveloping algebra [3]. This formula is derived from the Inönü–Wigner contraction

$$u(p, q) \rightsquigarrow Iu(p-1, q) \oplus \mathbb{R}, \quad (55)$$

and allows us to give a closed matrix expression for the invariants of the inhomogeneous algebras. The method can be enlarged to cover other contractions which are not necessarily inhomogeneous algebras. After adapting the formula to the special pseudo-unitary algebras $\mathfrak{su}(p, q)$, we analyse the problem of computing the invariants for the inhomogeneous special algebras $I\mathfrak{su}(p-1, q)$ applying the same procedure. However, the contraction of the invariants of $\mathfrak{su}(p, q)$ points out that the contraction of independent Casimir invariants does not necessarily lead to independent Casimir operators of the contraction. This fact is the consequence of altering the Cartan subalgebra in the contraction. To solve this problem, special pseudo-unitary algebras are contracted onto an extension $K\mathfrak{su}(p, q)$ of the inhomogeneous algebra $I\mathfrak{su}(p, q)$. The invariants W_k of this extension are shown to satisfy a functional equation (see (28)) that can be used to compute the invariants of $I\mathfrak{su}(p, q)$ by purely analytical methods. This procedure enlarges the extension method introduced in [13] for the special affine algebras $\mathfrak{sa}(n, \mathbb{R})$ and other algebras having only one Casimir operator.

Further, we have exhibited examples that show the validity of the argument when applied to other Lie algebras different from unitary algebras. In particular, matrix formulae for the Casimir operators of the inhomogeneous pseudo-orthogonal Lie algebras and two Inönü–Wigner contractions that naturally generalize the generalized Galilei and Carroll algebras have been given. The latter contractions of $I\mathfrak{so}(p+1, q)$ preserve the number of invariants [4]. Moreover, the structure of $C(p, q)$ shows that the procedure can be still valid for contractions of Lie algebras that are no more inhomogeneous.

Acknowledgments

The author expresses his gratitude to the referees for drawing the attention to some mistakes, as well as for valuable suggestions that helped to improve the manuscript. During the preparation of this work, the author was partially supported by the research project PR1/05-13283 of the UCM.

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